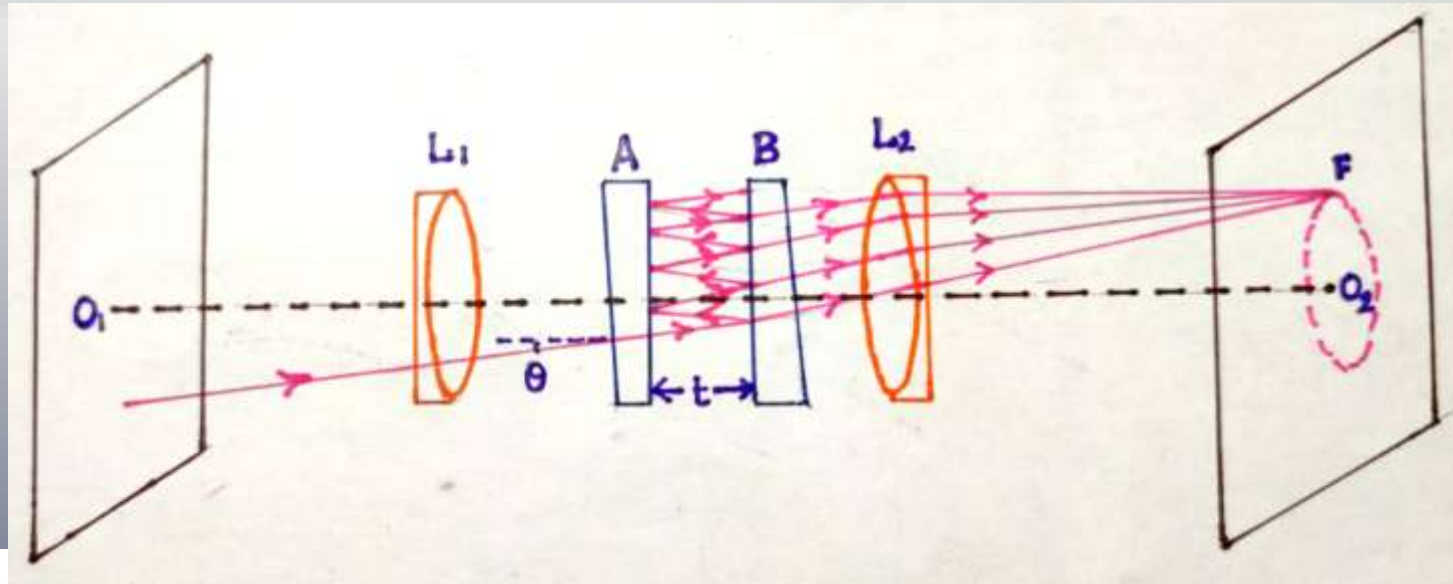


# FABRY– PEROT INTERFEROMETER

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## CONSTRUCTION

Fabry-Perot interferometer consists of two glass plates A and B separated by a distance. The inner surfaces of the plates are optically plane, exactly parallel and thin silvered so that about 70% of incident light gets reflected. The outer faces of the plates are also parallel to each other, but inclined to their respective inner faces.



If  $\theta$  is the angle of incidence on the silvered face of A, then, the path difference between successive rays is

$$2\mu t \cos \theta = 2 t \cos \theta \text{ as } \mu = 1 \text{ for air}$$

For bright fringe

$$2 t \cos \theta = n\lambda, n = 0, 1, 2, 3, 4, \dots$$

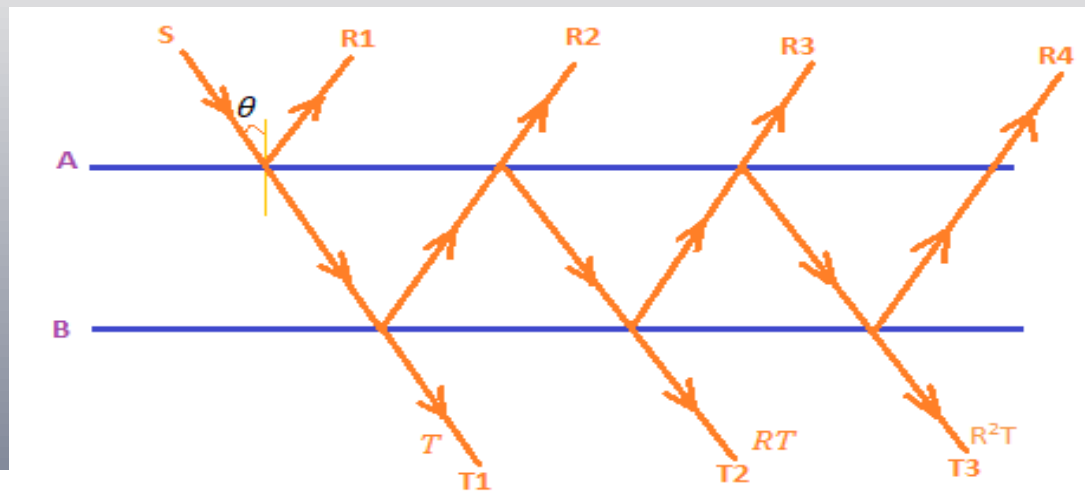
For all points passing through F with centre at  $O_2$  on the axis  $O_1 O_2$

When  $t$  is decreased, the ring shrinks and disappears at the centre. For a decrease of  $\lambda/2$ , one ring disappear at the centre.

Consider a plane wave of unit amplitude incident at an angle  $\theta$  on the glass plate.

Due to multiple reflections, a set of parallel reflected rays  $A_1R_1, A_2R_2, A_3R_3, \dots$ , and a set of parallel transmitted rays  $B_1T_1, B_2T_2, B_3T_3, \dots$  are produced.

The amplitudes of  $B_1T_1, B_2T_2, B_3T_3, \dots$  are  $T, RT, R^2T, R^3T, \dots$  respectively.



**Neglecting the small phase change due to reflection from silvered surfaces, the phase difference between two consecutive transmitted rays**

$$\delta = \frac{2\pi}{\lambda} \times 2t \cos\theta = \frac{4\pi t \cos\theta}{\lambda}$$

**Let the Incident wave be represented by**

$$y = a \sin \omega t = \sin \omega t \quad (\text{a is assumed to be unity})$$

**So, the transmitted rays can be represented as**

$$y_1 = T \sin \omega t$$

$$y_2 = RT \sin(\omega t - \delta)$$

$$y_3 = R^2 T \sin(\omega t - 2\delta), \text{ and so on}$$

**From the superposition principle**

$$Y = y_1 + y_2 + y_3 + \dots$$

**If A be the resultant amplitude of these rays and  $\phi$  be the phase difference, then we may write**

$$A \sin(\omega t - \phi) = T \sin \omega t + RT \sin(\omega t - \delta) + R^2 T \sin(\omega t - 2\delta) + \dots$$

$$A \sin \omega t \cos \phi - A \cos \omega t \sin \phi = T \sin \omega t + R T \sin \omega t \cos \delta - R T \cos \omega t \sin \delta + R^2 T \sin \omega t \cos 2\delta + R^2 T \cos \omega t \sin 2\delta + \dots$$

Equating the coefficients of  $\sin \omega t$  and  $\cos \omega t$  on both sides, we get

$$A \cos \phi = T + R T \cos \delta + R^2 T \cos 2\delta + \dots$$

$$A \sin \phi = R T \sin \delta + R^2 T \sin 2\delta + \dots$$

The resultant intensity

$$I = A^2 = (A \cos \phi + i A \sin \phi) \times (A \cos \phi - i A \sin \phi)$$

$$A \cos \phi + i A \sin \phi = T + R T (\cos \delta + i \sin \delta) + R^2 T (\cos 2\delta + i \sin 2\delta) + \dots$$

Which in terms of exponential can be put as

$$\begin{aligned} A \cos \phi + i A \sin \phi &= T + R T e^{i\delta} + R^2 T e^{2i\delta} + \dots \\ &= T(1 + R e^{i\delta} + R^2 e^{2i\delta} + \dots) \\ &= \frac{T}{1 - R e^{i\delta}} \end{aligned}$$

Similarly,

$$A \cos \phi - i A \sin \phi = \frac{T}{1 - R e^{-i\delta}}$$

The resultant intensity,

$$\begin{aligned} I = A^2 &= \frac{T}{1-Re^{i\delta}} \times \frac{T}{1-Re^{-i\delta}} \\ &= \frac{T^2}{1+R^2-2R \cos \delta} \\ &= \frac{T^2}{(1-R)^2+2R-2R \cos \delta} \\ &= \frac{T^2}{(1-R)^2+2R(1-\cos \delta)} \\ &= \frac{T^2}{(1-R)^2+4R \sin^2 \delta/2} \\ &= \frac{T^2}{(1-R)^2 \left[ 1 + \frac{4R}{(1-R)^2} \sin^2 \frac{\delta}{2} \right]} \end{aligned}$$

This expression is known as Airy's formula, shows that the resultant intensity depends upon the properties of the silver coating and  $\delta$ .